



$$B_{2g} \text{ orbital energies}$$

$$\Phi_{1} = \frac{1}{2} (\chi_{2} + \chi_{3} - \chi_{6} - \chi_{7}) \qquad \Phi_{2} = \frac{1}{2} (\chi_{1} + \chi_{4} - \chi_{5} - \chi_{8})$$

$$\begin{vmatrix} \alpha + \beta - E_{\mu} & \beta \\ \beta & \alpha - E_{\mu} \end{vmatrix} = 0 \implies E_{\mu}^{2} - \beta E_{\mu} - \beta^{2} = 0$$
(take $\alpha = 0$ - sets the energy zero)

$$E_{bonding} = \frac{\sqrt{5} + 1}{2} \beta; \quad E_{antibonding} = \frac{1 - \sqrt{5}}{2} \beta$$
MO coefficients? Plug each energy back into the secular eqn:
eg., for antibonding orb.:
$$\begin{bmatrix} \beta - \frac{1 - \sqrt{5}}{2} \beta & \beta \\ \beta & -\frac{1 - \sqrt{5}}{2} \beta \end{bmatrix} \begin{bmatrix} c_{1\mu} \\ c_{2\mu} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$B_{2g}$$
 orbital energies, cont.

For the antibonding orbital, the secular eqn. yields two eqns with the same solution:

$$\frac{c_2}{c_1} = -\frac{\sqrt{5}+1}{2}$$

Normalization of $\psi_{\text{antibonding}}$ demands $c_1^2 + c_2^2 = 1$

Solving:
$$c_1 = \sqrt{\frac{2}{5 + \sqrt{5}}} = 0.5257$$
; $c_2 = -0.8507$

Finally, we plug these into:
$$\psi_{\text{antibonding}} = c_1 \Phi_1 + c_2 \Phi_2$$

$$\Phi_{1} = \frac{1}{2} (\chi_{2} + \chi_{3} - \chi_{6} - \chi_{7}) \qquad \Phi_{2} = \frac{1}{2} (\chi_{1} + \chi_{4} - \chi_{5} - \chi_{8})$$

$$\psi_{\text{antibonding}} = 0.2629 (\chi_{2} + \chi_{3} - \chi_{6} - \chi_{7}) - 0.4253 (\chi_{1} + \chi_{4} - \chi_{5} - \chi_{8})$$









What is ESR (EPR)?

- ESR is the electron-spin analog of NMR.
- For a given applied field strength, ESR transitions occur at ~2000 times higher energy than NMR ($\mu_B/\mu_N = m_p/m_e = 1836.15$)
- Hyperfine splittings of ESR transitions arise from the interaction of the nuclear spins on electron spin transition energy.

H-atom Spin Hamiltonian $\mathcal{H}_{Spin} = g_e \mu_B \mathbf{H} \cdot \mathbf{S} - g_N \mu_N \mathbf{H} \cdot \mathbf{I} + a \mathbf{S} \cdot \mathbf{I}$ $a = \frac{8\pi}{3} g_e \mu_B g_N \mu_N |\psi(0)|^2$ • 1st term: electronic Zeeman • 2nd term: nuclear Zeeman • 3rd term: Fermi Contact hyperfine (isotropic) – magnitude depends on the electron density of the unpaired electron on the nucleus, $\psi(0)$.

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Hyperfine spitting in CH₃•

- Origin of proton hyperfine interaction: *Configuration interaction* must be accounted for, because at the orbital level π electrons have zero probability at the H-atom nuclei.

·H

 Get a quartet signal due to hyperfine splitting of three ¹H (spin-1/2) nuclei.

Methyl radical splitting diagram

• The four transitions for the methyl radical. + m_l states are lowest for $m_s = -1/_2$ and the - m_l state lowest for $m_s = 1/_2$, from the **I-S** term.



From Drago, "Physical Methods…" - Fig. 9-7 – corrected.

















Naphthalene MO Energies & Excited States

b _{3g} — -2.303	b _{3g} ——	b _{3g} —	b _{3g}
a _u — -1.618	a _u —	a _u —	a _u —
b _{1u} — -1.303 b _{3g} — -1.000 b _{2g} — -0.618	$b_{1u} - b_{3g} - b_{2g} + c_{3g} - b_{2g} - b_{2g} - b_{3g} - b$	$b_{1u} - b_{3g} + b_{2g} - b$	b_{1u} — b_{3g} — b_{2g} \uparrow — b_{2g} h — b_{2g}
	$\begin{array}{c} a_u \otimes b_{2g} = \\ B_{2u}(y) \end{array}$	$\begin{array}{c} a_u \otimes b_{3g} = \\ B_{3u}(x) \end{array}$	$b_{1u} \otimes b_{2g} = \mathbf{B}_{3u}(\mathbf{x})$
a _u 👉 0.618	au 🕁	au 🔶	au 🕂
b₁u 🕂 1.000	b1u 🕂	b1u 🕂	b1u 🔶
b _{3g} ₩ 1.303	b _{3g} 😽	b _{3g} 🛧	b _{3g} 😽
b _{2g} ₩ 1.618	b _{2g}	b _{2g}	b _{2g}
b _{1u} ₩ 2.303	b₁u ₩	b _{1u} 🔶	b _{1u} 🛧

0	bserved Tra	nsition Ener	rgies z	7_→ x
	Electronic	Transitions - Na	phthalene	
	E(cm ⁻¹)	Polarization	Assignment	
	31,800	Long axis (<i>x</i>)	$^1A_{1g} \rightarrow {}^1B_{3u}$	
	34,700	Short Axis (y)	$^1A_{1g} \rightarrow {}^1B_{2u}$	
	45,200	Long Axis (<i>x</i>)	$^1A_{1g} \rightarrow {}^1B_{3u}$	